

PHY509: SOLUTION HW #5.

P1. (a) $x = Q + s \cos \alpha$, $z = -s \sin \alpha$ and $X = Q$.

(b) $\delta x = \delta s \cos \alpha$, $\delta z = -\delta s \sin \alpha$ and $\delta X = 0$.

(c) $\delta x = \delta Q$, $\delta z = 0$ and $\delta X = \delta Q$.

(d) δQ does not change the gravitational potential of the system, therefore $Q_Q = -\partial V / \partial Q = 0$.

A finite δs changes the potential via δz and $Q_s = -\partial V / \partial s = -\partial(mgz) / \partial s = mg \sin \alpha$.

(e) $T = \frac{1}{2}m(\dot{x}^2 + \dot{z}^2) + \frac{1}{2}M\dot{X}^2 = \frac{1}{2}(m + M)\dot{Q}^2 + \frac{1}{2}m(\dot{s}^2 + 2\dot{Q}\dot{s} \cos \alpha)$.

Therefore the equations of motion for Q and s become, respectively,

$$(m + M)\ddot{Q} + m\ddot{s} \cos \alpha = 0 \quad (1)$$

$$m\ddot{s} + m\ddot{Q} \cos \alpha = mg \sin \alpha. \quad (2)$$

(f) The equation (1) is simply the conservation of total linear momentum. From this, we express \ddot{Q} in terms of \ddot{s} as $\ddot{Q} = -(m \cos \alpha / (m + M))\ddot{s}$. Substituting this into Eq. (2) we get

$$\ddot{s} - \frac{m \cos^2 \alpha}{m + M} \ddot{s} = g \sin \alpha, \quad \text{and} \quad \ddot{s} = \frac{(m + M)g \sin \alpha}{M + m \sin^2 \alpha}.$$

P2. (a) The length of the pendulum is $l = l_0 - \alpha t$. With the length l and the swing angle θ , $x = l \sin \theta$ and $z = -l \cos \theta$. Therefore the Lagrangian becomes

$$L = T - V = \frac{m}{2}[(\dot{l} \sin \theta + l\dot{\theta} \cos \theta)^2 + (\dot{l} \cos \theta - l\dot{\theta} \sin \theta)^2] + mgl \cos \theta = \frac{m}{2}(\alpha^2 + l^2\dot{\theta}^2) + mgl \cos \theta.$$

(b)

$$\frac{d}{dt} \left(\frac{\partial L}{\partial \dot{\theta}} \right) = \frac{d}{dt} (ml^2\dot{\theta}) = m(-2\alpha)l\dot{\theta} + ml^2\ddot{\theta}.$$

$$\text{Equation of motion becomes} \quad \ddot{\theta} - \frac{2\alpha}{l}\dot{\theta} + \frac{g}{l} \sin \theta = 0.$$

(c) Since the change in θ is much faster than the change in l , we approximate that l is a constant during a few oscillations. Also assume small oscillation. Then we propose a solution in the form $\theta = \theta_0 e^{-i\sigma t}$, which on substitution to the EOM gives an equation for σ , $\sigma^2 - i(2\alpha/l)\sigma - g/l = 0$. The solution for σ is

$$\sigma = i\frac{\alpha}{l} \pm \sqrt{\frac{g}{l} - \frac{\alpha^2}{l^2}}.$$

$$\text{Finally, } \theta(t) = \theta_0 \exp \left[\frac{\alpha t}{l_0 - \alpha t} \pm i \left(\frac{g}{l} - \frac{\alpha^2}{l^2} \right)^{1/2} t \right].$$

The amplitude increases with time, as can be inferred from experience.

Note: Since $\theta(t)$ increases with time, the small angle oscillation we used does not apply after certain time. But, at least, we get some qualitative ideas how it goes.